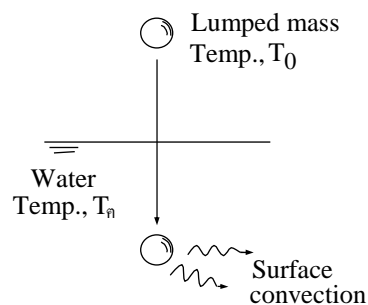


METHOD OF WEIGHTED RESIDUALS

METHOD OF WEIGHTED RESIDUALS

To understand the Method of Weighted Residuals (MWR), let's consider the example of transient thermal response of a lumped mass. Given,



A = surface area
 V = volume
 ρ = density
 c = specific heat
 h = convection coefficient
 τ = time
 T = Temperature

Want: Lumped mass temperature, $T = T(\tau) = ?$

TRANSIENT RESPONSE OF LUMPED MASS

Conservation of energy:

Energy in - Energy out = Energy stored

$$0 - h A (T - T_{\infty}) = \rho c V \frac{\partial T}{\partial \tau}$$

$$\rho c V \frac{\partial T}{\partial \tau} + h A (T - T_{\infty}) = 0$$

with initial condition: $T(\tau = 0) = T_0$

To easily solve this ODE, let's denote,

$$\text{Non-dimensional temperature } \bar{x} = \frac{T - T_{\infty}}{T_0 - T_{\infty}}$$

$$\text{Non-dimensional time } t = \frac{h A \tau}{\rho c V}$$

then the ODE & IC become,

TRANSIENT RESPONSE OF LUMPED MASS

ODE: $\frac{d\bar{x}}{dt} + \bar{x} = 0$

IC: $\bar{x}(t = 0) = 1$

To solve for exact solution, $\bar{x}(t)$, we separate the variables,

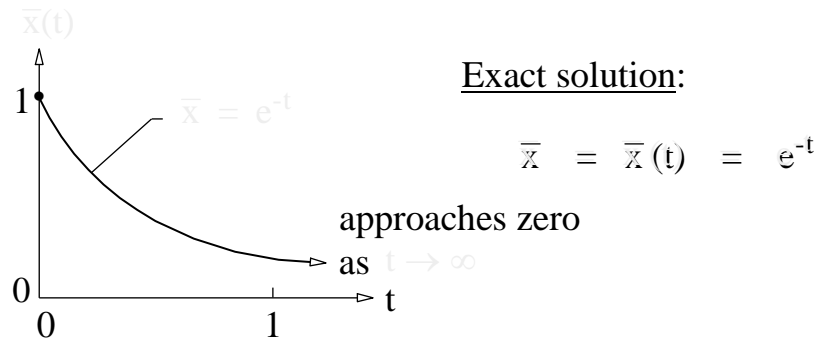
$$\frac{d\bar{x}}{\bar{x}} = -dt$$

then integrate $\ln \bar{x} = -t + B$

where B is the integrating constant which is zero after applying the initial condition. Thus, the exact solution is,

$$\bar{x} = \bar{x}(t) = e^{-t}$$

TRANSIENT RESPONSE OF LUMPED MASS



Now, if we don't know the exact solution, and we want to use the method of weighted residuals to find approximate solution for $0 < t < 1$.

METHOD OF WEIGHTED RESIDUALS

Assume an approximate solution in the form,

$$x(t) = 1 + C_1 t + C_2 t^2$$

where C_1 and C_2 are constants to be determined. Note that if we substitute this approx. sol. into the diff. eq.

$$\frac{dx}{dt} + x \neq 0 \quad \text{in fact} \quad = R$$

i.e., $R = (0 + C_1 + 2C_2 t) + (1 + C_1 t + C_2 t^2)$

or, $R = R(t) = 1 + C_1(1 + t) + C_2(2t + t^2)$

So, the idea is to determine C_1 and C_2 such that the Residual $R \rightarrow 0$.

METHOD OF WEIGHTED RESIDUALS

Four approaches:

- (1) Point collocation Set residuals at selected points to be zero.
- (2) Subdomain collocation Set areas under the residual curve to be zero.
- (3) Galerkin Include weighting functions. Widely used in FE method.
- (4) Least squares Square the residual with minimization process.

POINT COLLOCATION

Set the residuals at selected points to be zero, i.e.,

$$R(t_i) = 0 \quad i = 1, 2, \dots, \text{no. of unknowns}$$

Since there are two unknowns of C_1 and C_2 , thus we need 2 eqs. of,

$$R(t_1) = 0 \quad \text{and} \quad R(t_2) = 0$$

where t_1 and t_2 should be in the range of consider 0 to 1.

E.g., select,

$$t_1 = \frac{1}{3}; \quad R\left(\frac{1}{3}\right) = 1 + C_1\left(1 + \frac{1}{3}\right) + C_2\left(2\left(\frac{1}{3}\right) + \left(\frac{1}{3}\right)^2\right) = 0$$

$$t_2 = \frac{2}{3}; \quad R\left(\frac{2}{3}\right) = 1 + C_1\left(1 + \frac{2}{3}\right) + C_2\left(2\left(\frac{2}{3}\right) + \left(\frac{2}{3}\right)^2\right) = 0$$

Solve to get, $C_1 = -0.9310$ and $C_2 = 0.3103$

Then, $x(t) = 1 - 0.9310 t + 0.3103 t^2$

SUBDOMAIN COLLOCATION

Set areas under the residual curve to be zero.

Since there are 2 unknowns, thus we need 2 eqs., e.g.

$$\int_0^{t_1} R(t) dt = 0 \quad \text{and} \quad \int_{t_1}^1 R(t) dt = 0$$

If we select $t_1 = \frac{1}{2}$, then,

$$\int_0^{1/2} R(t) dt; \quad \frac{1}{2} + \frac{5}{8}C_1 + \frac{7}{24}C_2 = 0$$

$$\int_{1/2}^1 R(t) dt; \quad \frac{1}{2} + \frac{7}{8}C_1 + \frac{25}{24}C_2 = 0$$

Solve to get, $C_1 = -0.9474$ and $C_2 = 0.3158$

Then, the approximate solution is,

$$x(t) = 1 - 0.9474 t + 0.3158 t^2$$

GALERKIN APPROACH

Multiply the residual by some weighting functions, perform integration over the entire domain, and set to be zero.

$$\int_0^1 R(t) W_i(t) dt = 0 \quad i = 1, 2, \dots, \text{No. of unknowns}$$

where $W_i(t)$ is the weighting functions which are normally selected as the terms that weight with the unknowns, i.e.,

$$x(t) = 1 + C_1 \underbrace{t}_{W_1} + C_2 \underbrace{t^2}_{W_2}$$

here, we use $W_1 = t$ and $W_2 = t^2$, then substitute into the above eq. to get,

GALERKIN APPROACH

$$\int_0^1 R(t) t dt = 0; \quad \frac{1}{2} + \frac{5}{6}C_1 + \frac{11}{12}C_2 = 0$$

$$\int_0^1 R(t) t^2 dt = 0; \quad \frac{1}{3} + \frac{7}{12}C_1 + \frac{7}{10}C_2 = 0$$

Solve to get, $C_1 = -0.9143$ and $C_2 = 0.2857$

Thus, the approximate solution is,

$$x(t) = 1 - 0.9143 t + 0.2857 t^2$$

Note: We will use this approach to derive FE eqs. later. Since distribution of element variable is,

$$\phi = \underbrace{N_1}_{\bar{W}_1} \phi_1 + \underbrace{N_2}_{\bar{W}_2} \phi_2$$

Here, the nodal unknowns ϕ_1 and ϕ_2 are similar to C_1 and C_2 . The element interpolation functions N_1 and N_2 will be selected as the weighting functions.

LEAST SQUARES APPROACH

Square the residual, perform integration over entire domain, and minimize with respect to the unknowns, i.e.,

$$\frac{\partial}{\partial C_i} \int_0^1 R^2(t) dt = 0 \quad i = 1, 2, \dots, \text{No. of unknowns}$$

Since there are 2 unknowns, thus we have 2 eqs.,

$$\frac{\partial}{\partial C_1} \int_0^1 R^2(t) dt = 0; \quad \frac{3}{2} + \frac{7}{3}C_1 + \frac{9}{4}C_2 = 0$$

$$\frac{\partial}{\partial C_2} \int_0^1 R^2(t) dt = 0; \quad \frac{4}{3} + \frac{9}{4}C_1 + \frac{33}{15}C_2 = 0$$

Solve to get, $C_1 = -0.9427$ and $C_2 = 0.3110$

Then, the approximate solution is,

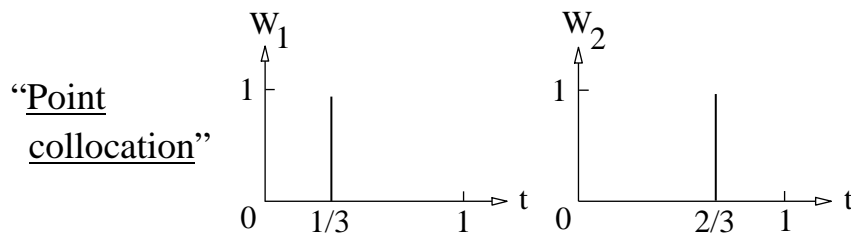
$$x(t) = 1 - 0.9427 t + 0.3110 t^2$$

METHOD OF WEIGHTED RESIDUALS

In conclusion, these four approaches can be written in the form,

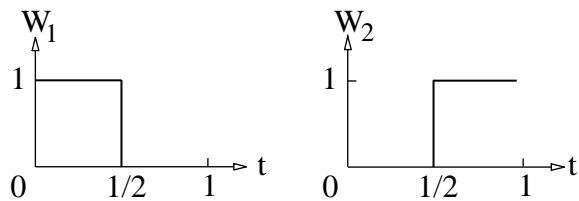
$$\int_0^1 R(t) W_i(t) dt = 0$$

where W_i are the weighting functions that can be described as follows,

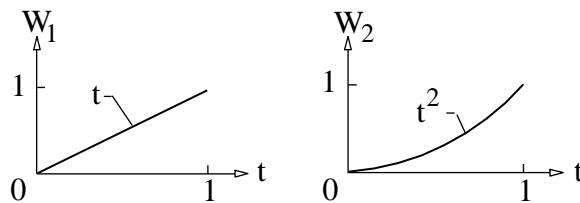


METHOD OF WEIGHTED RESIDUALS

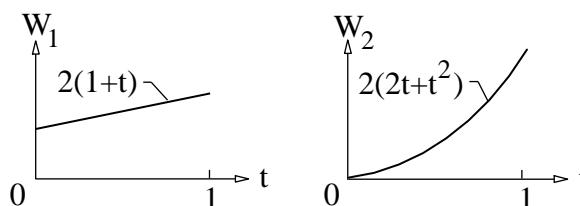
“Subdomain collocation”



“Galerkin”

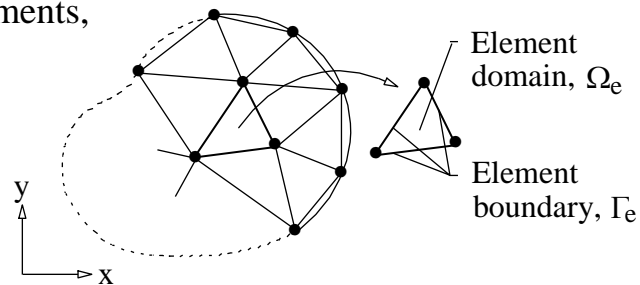


“Least squares”



GENERAL PROCEDURE OF MWR

Step 1: Discretize the given domain into a number of elements,



Then identify the governing differential eq(s) of the problem,

$$L(\bar{\phi}) = 0$$

where L denotes the differential operator and $\bar{\phi}$ is the exact solution.

GENERAL PROCEDURE OF MWR

Step 2: Assume element approximate solution in the form,

$$\phi = \phi(x, y) = \sum_{i=1}^m N_i \phi_i = \begin{matrix} [N] \\ (1 \times m) \end{matrix} \begin{matrix} \{\phi\} \\ (1 \times m) \end{matrix}$$

where $\phi = \phi(x, y)$ = element approximate solution
 $N_i = N_i(x, y)$ = element interpolation functions
 ϕ_i = element nodal unknowns
 m = number of element nodes

Step 3: Apply the method of weighted residuals for a typical element. If we substitute the approximate solution, $\phi(x, y)$, into the governing diff. eq.,

$$L(\phi) \neq 0 \text{ in general, but } L(\phi) = \text{Error} = R$$

GENERAL PROCEDURE OF MWR

Here R is called the Residual, i.e.,

$$R = L(\phi) = L(\{N\}\{\phi\}) = L\left(\sum_{i=1}^m N_i \phi_i\right)$$

Then apply the method of weighted residuals (Galerkin approach):

Multiply the Residual R by some weighting functions, perform integration over element domain, and set to be zero.

$$\int_{\Omega^{(e)}} W_i R \, d\Omega = 0 \quad i = 1, 2, \dots, m$$

Note: If $W_i = N_i$ \Rightarrow called “Bubnov-Galerkin” approach
 $W_i \neq N_i$ \Rightarrow called “Petrov-Galerkin” approach

GENERAL PROCEDURE OF MWR

Step 4: Perform integration by parts (or use Gauss’s theorem).

$$\begin{aligned} \int_{\Omega^{(e)}} W_i R \, d\Omega &= \int_{\Omega^{(e)}} W_i L\left(\sum_{i=1}^m N_i \phi_i\right) d\Omega \\ &= \underbrace{\int_{\Omega^{(e)}} (W_i, N_i, \phi_i) d\Omega}_{\text{Associated with element domain, } \Omega^{(e)}} + \underbrace{\int_{\Gamma^{(e)}} (W_i, N_i, \phi_i) d\Gamma}_{\text{Associated with element boundary, } \Gamma^{(e)}} = 0 \end{aligned}$$

Step 5: Apply boundary conditions on the element boundary integral term $\int_{\Gamma^{(e)}}$ as needed.

GENERAL PROCEDURE OF MWR

Step 6: Write the element eqs. in matrix form,

$$\begin{array}{ccc} [\mathbf{K}] & \{\phi\} & = & \{F\} \\ \text{(mxm)} & \text{(mx1)} & & \text{(mx1)} \end{array}$$

where $[\mathbf{K}]$ = Element stiffness matrix

$\{\phi\}$ = Vector of element nodal unknowns

$\{F\}$ = Vector of element nodal loads

Then assemble all element eqs., apply all BC's on the system eqs., and solve for nodal unknowns.

METHOD OF WEIGHTED RESIDUALS

Objective: Derive finite element equations from governing differential equation.

Simplest example is 1-D Poisson's eq.:

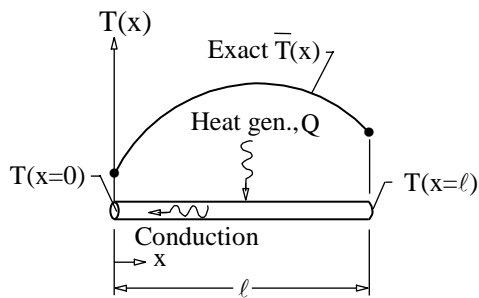
$$\frac{d^2 u}{dx^2} = f(x)$$

where $u = u(x)$. Such differential eq. may represent the problems of,

- Bar with its own weight
- 1-D viscous flow
- 1-D heat conduction with internal heat generation

etc.

CONDUCTION WITH HEAT GENERATION



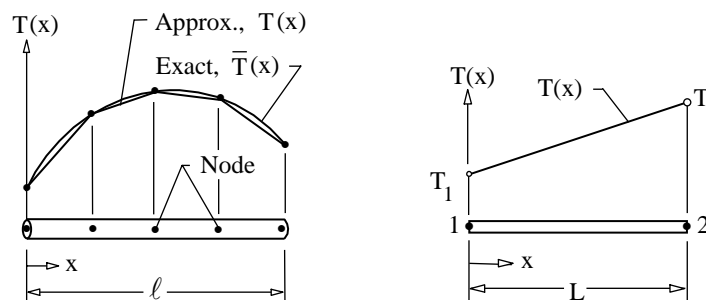
Differential equation:

$$kA \frac{d^2 \bar{T}}{dx^2} = -QA$$

where k is thermal conductivity and A is cross-sectional area.

Boundary conditions: $T(x=0) = T_0$ and $T(x=l) = T_\ell$

DISCRETIZATION & INTERPOLATION FUNC.



Assume linear element temperature distribution,

$$T(x) = ax + b$$

$$\text{At } x = 0; \quad T(x=0) = T_1 = b$$

$$\text{At } x = L; \quad T(x=L) = T_2 = aL + b$$

$$\text{Then } a = \frac{T_2 - T_1}{L} \quad \text{and} \quad b = T_1$$

ELEMENT INTERPOATION FUNCTIONS

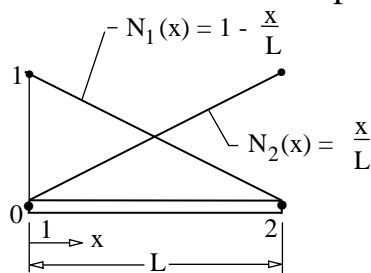
Then
$$T(x) = \left(1 - \frac{x}{L}\right)T_1 + \left(\frac{x}{L}\right)T_2$$

which can be written in the form,

$$T(x) = N_1 T_1 + N_2 T_2 = \begin{bmatrix} N_1 & N_2 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} = \begin{bmatrix} N \end{bmatrix} \begin{Bmatrix} T \end{Bmatrix}$$

(1x2) (2x1)

where the element interpolation functions,



$$N_1 = 1 - \frac{x}{L}$$

$$N_2 = \frac{x}{L}$$

with properties of,

$$N_i = \begin{cases} 1 & \text{at node } i \\ 0 & \text{at other node} \end{cases}$$

METHOD OF WEIGHTED RESIDUALS

If we can solve for exact solution, \bar{T} , and substitute into the differential eq., then RHS will be identically zero, i.e.,

$$kA \frac{d^2 \bar{T}}{dx^2} + QA = 0$$

However, in general, we don't know exact solution. Thus, if we substitute the approximate solution, T , RHS will not be zero, but will be equal to a Residual, R :

$$kA \frac{d^2 T}{dx^2} + QA = R$$

The idea is thus to minimize R so that error is minimum.

METHOD OF WEIGHTED RESIDUALS

The method of weighted residuals is to multiply R by some weighting functions W_i , perform integration over element domain, and set to zero, i.e.:

$$\int_0^L W_i R \, dx = 0$$

Since the element has two unknowns, thus we need 2 eqs., i.e., $i = 1, 2$. Substitute the residual to get,

$$\int_0^L W_i \left(kA \frac{d^2 T}{dx^2} + QA \right) dx = 0$$

Expand,

$$\int_0^L W_i kA \frac{d^2 T}{dx^2} dx + \int_0^L W_i QA \, dx = 0 \quad i = 1, 2$$

METHOD OF WEIGHTED RESIDUALS

Apply integration by parts to the first integral term in order to produce the boundary term, i.e.,

$$\int_0^L \underbrace{W_i}_u \underbrace{kA \frac{d^2 T}{dx^2}}_{dv} dx = W_i kA \frac{dT}{dx} \Big|_0^L - \int_0^L kA \frac{dT}{dx} \frac{dW_i}{dx} dx$$

by using the formula, $\int_0^L u \, dv = u v \Big|_0^L - \int_0^L v \, du$

where $u = W_i$ then $du = \frac{dW_i}{dx} dx$

and $dv = kA \frac{d^2 T}{dx^2} dx$ then $v = kA \frac{dT}{dx}$

Then, the element eqs. becomes,

METHOD OF WEIGHTED RESIDUALS

$$\int_0^L k A \frac{dW_i}{dx} \frac{dT}{dx} dx = W_i k A \frac{dT}{dx} \Big|_0^L + \int_0^L W_i Q A dx$$

Recall $T = \begin{bmatrix} N_1 & N_2 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$ then $\frac{dT}{dx} = \begin{bmatrix} \frac{dN_1}{dx} & \frac{dN_2}{dx} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$

Substitute and write out for $i = 1, 2$ to get 2 equations as,

$$i = 1; \int_0^L k A \frac{dW_1}{dx} \begin{bmatrix} \frac{dN_1}{dx} & \frac{dN_2}{dx} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} dx = W_1 k A \frac{dT}{dx} \Big|_0^L + \int_0^L W_1 Q A dx$$

$$i = 2; \int_0^L k A \frac{dW_2}{dx} \begin{bmatrix} \frac{dN_1}{dx} & \frac{dN_2}{dx} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} dx = W_2 k A \frac{dT}{dx} \Big|_0^L + \int_0^L W_2 Q A dx$$

METHOD OF WEIGHTED RESIDUALS

For standard FE., we select $W_i = N_i$ which is known as Bubnov-Galerkin approach, we will get element equations in the form,

$$\begin{aligned} & \int_0^L k A \begin{Bmatrix} \frac{dN_1}{dx} \\ \frac{dN_2}{dx} \end{Bmatrix} \begin{bmatrix} \frac{dN_1}{dx} & \frac{dN_2}{dx} \end{bmatrix} dx \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} \\ & = \left(\begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} k A \frac{dT}{dx} \right) \Big|_0^L + \int_0^L \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} Q A dx \end{aligned}$$

METHOD OF WEIGHTED RESIDUALS

Or, in short: $[K_C]\{T\} = \{Q_C\} + \{Q_Q\}$

where $[K_C]$ = Element conduction matrix

$\{T\}$ = Vector of element nodal temperatures

$\{Q_C\}$ = Vector of element nodal heat fluxes

$\{Q_Q\}$ = Vector of element nodal heat generations

METHOD OF WEIGHTED RESIDUALS

Conduction matrix $[K_C] = \int_0^L k A \begin{Bmatrix} \frac{dN_1}{dx} \\ \frac{dN_2}{dx} \end{Bmatrix} \begin{bmatrix} \frac{dN_1}{dx} & \frac{dN_2}{dx} \end{bmatrix} dx$

Since $N_1 = 1 - \frac{x}{L}$ and $N_2 = \frac{x}{L}$
 then $\frac{dN_1}{dx} = -\frac{1}{L}$ and $\frac{dN_2}{dx} = \frac{1}{L}$

Substitute $[K_C] = \int_0^L k A \begin{Bmatrix} -\frac{1}{L} \\ \frac{1}{L} \end{Bmatrix} \begin{bmatrix} -\frac{1}{L} & \frac{1}{L} \end{bmatrix} dx = \int_0^L \frac{k A}{L^2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} dx$

If thermal conductivity k & cross-sectional area A are constant, then

$$[K_C] = \frac{k A}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

ELEMENT LOAD VECTORS

Load vector due to internal heat generation

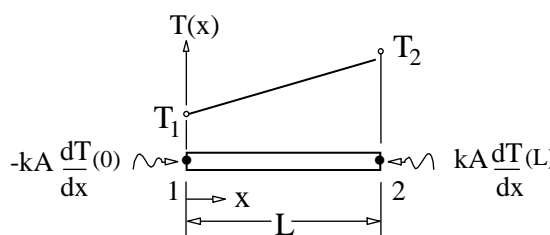
$$\{Q_Q\} = \int_0^L \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} Q A dx = \int_0^L \begin{Bmatrix} 1 - \frac{x}{L} \\ \frac{x}{L} \end{Bmatrix} Q A dx = Q A L \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

Load vector due to nodal heat fluxes

$$\begin{aligned} \{Q_C\} &= \left(\begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} k A \frac{dT}{dx} \right) \Big|_0^L = \begin{Bmatrix} \left(N_1 k A \frac{dT}{dx} \right) \Big|_0^L \\ \left(N_2 k A \frac{dT}{dx} \right) \Big|_0^L \end{Bmatrix} \\ &= \begin{Bmatrix} N_1(L) k A \frac{dT}{dx}(L) - N_1(0) k A \frac{dT}{dx}(0) \\ N_2(L) k A \frac{dT}{dx}(L) - N_2(0) k A \frac{dT}{dx}(0) \end{Bmatrix} \end{aligned}$$

But $N_1(0) = 1, N_1(L) = 0, N_2(0) = 0, N_2(L) = 1$, thus

ELEMENT LOAD VECTORS



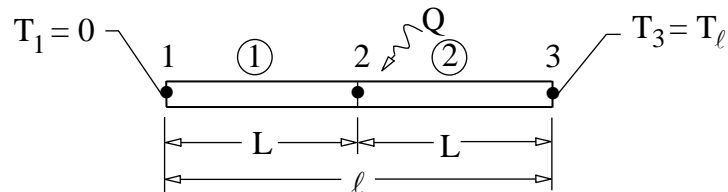
$$\{Q_C\} = \begin{Bmatrix} -k A \frac{dT}{dx}(0) \\ k A \frac{dT}{dx}(L) \end{Bmatrix}$$

Conclusion of element equations:

$$\begin{aligned} [K_C] \{T\} &= \{Q_C\} + \{Q_Q\} \\ \frac{k A}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} &= \begin{Bmatrix} -k A \frac{dT}{dx}(0) \\ k A \frac{dT}{dx}(L) \end{Bmatrix} + Q A L \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix} \end{aligned}$$

METHOD OF WEIGHTED RESIDUALS

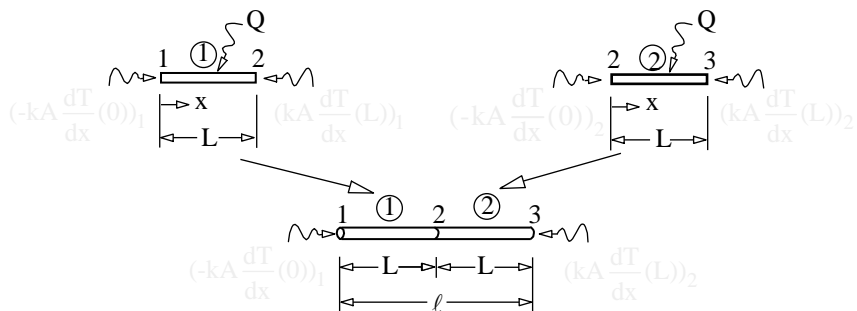
Example A rod divided into 2 elements with specified temp. at both ends as shown is subjected to internal heat generation Q . Compute the temperature at node ②.



Typical element eqs., e.g., for element no. ①:

$$\frac{kA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} = \begin{Bmatrix} (-kA \frac{dT}{dx}(0))_1 \\ (kA \frac{dT}{dx}(L))_1 \end{Bmatrix} + QAL \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

ASSEMBLING OF ELEMENT EQUATIONS



$$\frac{kA}{L} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1+1 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \end{Bmatrix} = \begin{Bmatrix} (-kA \frac{dT}{dz}(0))_1 \\ (kA \frac{dT}{dz}(L))_1 + (-kA \frac{dT}{dz}(0))_2 \\ (kA \frac{dT}{dz}(L))_2 \end{Bmatrix} + QAL \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} + \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

SYSTEM OF EQUATIONS

Heat fluxes at node 2 must be continuous, and apply BC's:

$$\frac{kA}{L} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 = 0 \\ T_2 = ? \\ T_3 = T_\ell \end{Bmatrix} = \begin{Bmatrix} (-kA \frac{dT}{dx}(0))_1 \\ 0 \\ (kA \frac{dT}{dx}(L))_2 \end{Bmatrix} + QAL \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

Then use second eq. to solve for temperature at node 2,

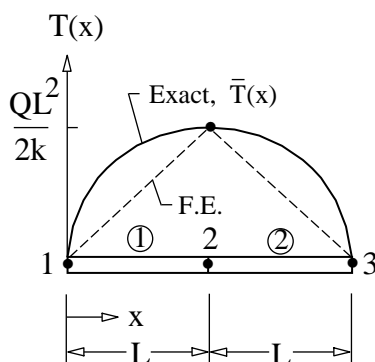
$$\begin{aligned} \frac{kA}{L} (0 + 2T_2 - T_\ell) &= 0 + QAL \\ T_2 &= \frac{QL^2}{2k} + \frac{T_\ell}{2} \end{aligned}$$

Note: Identical procedure can be used for more elements.

USE OF SOLUTION SYMMETRY

If $T_\ell = 0$, then the temperature at node 2 becomes,

$$T_2 = \frac{QL^2}{2k}$$



Note that the exact temperature distribution in this case is,

$$\bar{T}(x) = \frac{QL^2}{2k} \left(2\frac{x}{L} - \frac{x^2}{L^2} \right)$$

which is symmetry.

USE OF SOLUTION SYMMETRY

Due to symmetry of solution, only half of the rod can be used in modeling. Here, if use 1 element,

$$\frac{kA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} = \begin{Bmatrix} (-kA \frac{dT}{dx}(0))_1 \\ (kA \frac{dT}{dx}(L))_1 \end{Bmatrix} + QAL \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

and apply the condition of no heat flow across node 2, i.e.,

$$\left(kA \frac{dT}{dx}(L) \right)_1 = 0$$

Then the element eqs. become,

$$\frac{kA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 = 0 \\ T_2 = ? \end{Bmatrix} = \begin{Bmatrix} (-kA \frac{dT}{dx}(0))_1 \\ 0 \end{Bmatrix} + QAL \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

USE OF SOLUTION SYMMETRY

Then use the second eq. to solve for T_2 ,

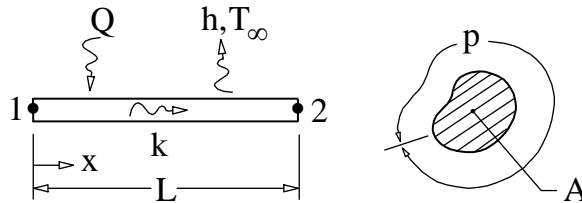
$$\begin{aligned} \frac{kA}{L} (0 + T_2) &= 0 + \frac{QAL}{2} \\ T_2 &= \frac{QL^2}{2k} \end{aligned}$$

Remark: For practical problems in 3-D, always take advantage of solution symmetry in order to reduce total number of elements and unknowns.

<u>Problem</u>	<u>Reduction factor</u>
1-D	2
2-D	4
3-D	8

ROD WITH SURFACE CONVECTION

If the rod has additional surface convection,



The corresponding differential equation is,

$$k A \frac{d^2 T}{dx^2} - h p (T - T_\infty) + Q A = 0$$

where h = convection coefficient
 p = perimeter
 T_∞ = surrounding medium temperature

ROD WITH SURFACE CONVECTION

Apply MWR,

$$\int_0^L W_i R dx = 0 \quad i = 1, 2$$

$$\int_0^L W_i \left(k A \frac{d^2 T}{dx^2} - h p (T - T_\infty) + Q A \right) dx = 0$$

$$\int_0^L W_i k A \frac{d^2 T}{dx^2} dx + \int_0^L W_i Q A dx$$

$$- \int_0^L W_i h p T dx + \int_0^L W_i h p T_\infty dx = 0$$

with $W_i = N_i$ and $T = [N] \{T\}$, the last two integral terms become,

ROD WITH SURFACE CONVECTION

$$\int_0^L W_i h p T dx \Rightarrow \int_0^L h p \{N\} [N] dx \{T\}$$

$$= [K_h] \{T\}$$

$$\int_0^L W_i h p T_\infty dx \Rightarrow \int_0^L h p T_\infty \{N\} dx$$

$$= \{Q_h\}$$

where $[K_h]$ = Convection matrix

$\{Q_h\}$ = Convection load vector

ROD WITH SURFACE CONVECTION

The finite element eqs. are in the form,

$$[K_c] \{T\} + [K_h] \{T\} = \{Q_c\} + \{Q_Q\} + \{Q_h\}$$

or, in detail,

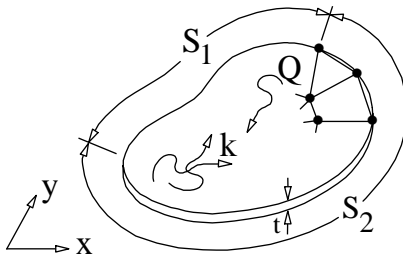
$$\frac{k A}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix} + h p L \begin{bmatrix} \frac{1}{3} & \frac{1}{6} \\ \frac{1}{6} & \frac{1}{3} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \end{Bmatrix}$$

$$= \begin{Bmatrix} -k A \frac{dT}{dx}(0) \\ k A \frac{dT}{dx}(L) \end{Bmatrix} + Q A L \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix} + h p L T_\infty \begin{Bmatrix} \frac{1}{2} \\ \frac{1}{2} \end{Bmatrix}$$

which can be used in many applications, such as heat transfer in motor fins, etc.

TWO-DIMENSIONAL HEAT TRANSFER

Plate with Internal Heat Generation



Governing differential equation:

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) = -Q$$

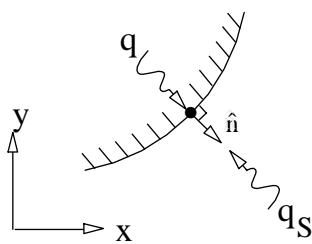
Boundary conditions:

(1) Specified temperature along edge S_1 :

$$T(x, y) = T_1(x, y)$$

TWO-DIMENSIONAL HEAT TRANSFER

(2) Specified surface heating along edge S_2 . From Fourier's law,



$$q = -k \frac{\partial T}{\partial x} n_x - k \frac{\partial T}{\partial y} n_y$$

where n_x and n_y are direction cosines of unit vector \hat{n} normal to the edge. For a specified heating q_s into edge, then

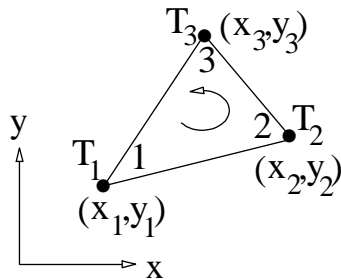
$$q_s = -q = k \frac{\partial T}{\partial x} n_x + k \frac{\partial T}{\partial y} n_y$$

This means, for insulated edge,

$$k \frac{\partial T}{\partial x} n_x + k \frac{\partial T}{\partial y} n_y = 0$$

TRIANGULAR ELEMENT

Element Interpolation Functions



Assume a flat temperature distribution over the element:

$T(x, y) = \alpha_1 + \alpha_2 x + \alpha_3 y$
 where α_i , $i = 1, 2, 3$, are constants to be determined from conditions at nodes,

At node 1: $T(x_1, y_1) = T_1 = \alpha_1 + \alpha_2 x_1 + \alpha_3 y_1$

At node 2: $T(x_2, y_2) = T_2 = \alpha_1 + \alpha_2 x_2 + \alpha_3 y_2$

At node 3: $T(x_3, y_3) = T_3 = \alpha_1 + \alpha_2 x_3 + \alpha_3 y_3$

Solve for α_i , $i = 1, 2, 3$, and rearrange terms to get,

TRIANGULAR ELEMENT

$$T(x, y) = \begin{bmatrix} N_1(x, y) & N_2(x, y) & N_3(x, y) \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \end{Bmatrix} = \begin{bmatrix} \mathbf{N} \end{bmatrix} \{ \mathbf{T} \}$$

(1x3)(3x1)

where $\{T\}$ is the vector of nodal temperatures, and $[N]$ is the element interpolation matrix defined by,

$$N_i(x, y) = \frac{1}{2A} (a_i + b_i x + c_i y) \quad i = 1, 2, 3$$

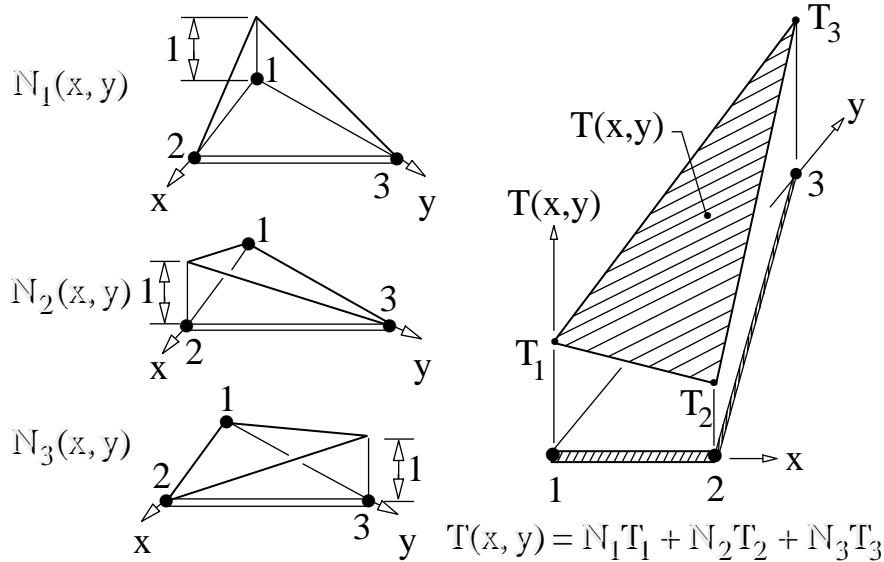
Here $A = \text{Area of triangle}$

$$= \frac{1}{2} [x_2(y_3 - y_1) + x_1(y_2 - y_3) + x_3(y_1 - y_2)]$$

and

$$\begin{aligned} a_1 &= x_2 y_3 - x_3 y_2 & b_1 &= y_2 - y_3 & c_1 &= x_3 - x_2 \\ a_2 &= x_3 y_1 - x_1 y_3 & b_2 &= y_3 - y_1 & c_2 &= x_1 - x_3 \\ a_3 &= x_1 y_2 - x_2 y_1 & b_3 &= y_1 - y_2 & c_3 &= x_2 - x_1 \end{aligned}$$

TRIANGULAR ELEMENT INTERPOLATIONS



DERIVATION OF ELEMENT EQUATIONS

If \bar{T} is the exact temperature solution of the plate, then

$$\frac{\partial}{\partial x} \left(k \frac{\partial \bar{T}}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial \bar{T}}{\partial y} \right) + Q = 0$$

Since we use approximate T , then RHS of the above equation is not zero, but is equal to a Residual R , i.e.

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + Q = R$$

Apply MWR: $\int_{\Omega^{(e)}} W_i R \, d\Omega = 0$

where $\Omega^{(e)}$ is the element volume. Substitute,

$$\int_{\Omega^{(e)}} W_i \left(\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) + Q \right) d\Omega = 0$$

DERIVATION OF ELEMENT EQUATIONS

Expand,

$$\int_{\Omega^{(e)}} W_i \left(\frac{\partial}{\partial x} \left(k \frac{\partial \Gamma}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial \Gamma}{\partial y} \right) \right) d\Omega + \int_{\Omega^{(e)}} W_i Q d\Omega = 0$$

As in 1-D, we apply integration by parts to the first term,
Here is to use Gauss's theorem,

$$\int_{\Omega^{(e)}} u (\nabla \cdot \vec{V}) d\Omega = \int_{\Gamma^{(e)}} u (\vec{V} \cdot \hat{n}) d\Gamma - \int_{\Omega^{(e)}} (\nabla u \cdot \vec{V}) d\Omega$$

Let $u = W_i$

$$\left. \begin{aligned} \nabla &= \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} \\ \vec{V} &= k \frac{\partial \Gamma}{\partial x} \hat{i} + k \frac{\partial \Gamma}{\partial y} \hat{j} \end{aligned} \right\} (\nabla \cdot \vec{V}) = \left(\frac{\partial}{\partial x} \left(k \frac{\partial \Gamma}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial \Gamma}{\partial y} \right) \right)$$

DERIVATION OF ELEMENT EQUATIONS

Then the terms on RHS of Gauss's theorem which states
as,

$$\int_{\Omega^{(e)}} u (\nabla \cdot \vec{V}) d\Omega = \int_{\Gamma^{(e)}} u (\vec{V} \cdot \hat{n}) d\Gamma - \int_{\Omega^{(e)}} (\nabla u \cdot \vec{V}) d\Omega$$

using $\hat{n} = n_x \hat{i} + n_y \hat{j}$, are:

$$\begin{aligned} \vec{V} \cdot \hat{n} &= k \frac{\partial \Gamma}{\partial x} n_x + k \frac{\partial \Gamma}{\partial y} n_y \\ u(\vec{V} \cdot \hat{n}) &= W_i \left(k \frac{\partial \Gamma}{\partial x} n_x + k \frac{\partial \Gamma}{\partial y} n_y \right) \\ \nabla u &= \frac{\partial W_i}{\partial x} \hat{i} + \frac{\partial W_i}{\partial y} \hat{j} \\ \nabla u \cdot \vec{V} &= \frac{\partial W_i}{\partial x} k \frac{\partial \Gamma}{\partial x} + \frac{\partial W_i}{\partial y} k \frac{\partial \Gamma}{\partial y} \end{aligned}$$

DERIVATION OF ELEMENT EQUATIONS

Thus the element equations with $W_i = N_i$ are,

$$\int_{\Gamma^{(e)}} N_i \left(k \frac{\partial T}{\partial x} n_x + k \frac{\partial T}{\partial y} n_y \right) d\Gamma - \int_{\Omega^{(e)}} \left(\frac{\partial N_i}{\partial x} k \frac{\partial T}{\partial x} + \frac{\partial N_i}{\partial y} k \frac{\partial T}{\partial y} \right) d\Omega + \int_{\Omega^{(e)}} N_i Q d\Omega = 0 \quad i = 1, 2, 3$$

Imposing the boundary conditions along edges, the element equations can be written in matrix form (total of 3 eqs.) as,

$$\int_{\Omega^{(e)}} \left(\left\{ \frac{\partial N}{\partial x} \right\} k \frac{\partial T}{\partial x} + \left\{ \frac{\partial N}{\partial y} \right\} k \frac{\partial T}{\partial y} \right) d\Omega = \int_{\Gamma^{(e)}} \{N\} \left(k \frac{\partial T}{\partial x} n_x + k \frac{\partial T}{\partial y} n_y \right) d\Gamma + \int_{\Omega^{(e)}} \{N\} Q d\Omega + \int_{\Gamma^{(e)}} \{N\} q_s d\Gamma$$

DERIVATION OF ELEMENT EQUATIONS

Since we assume element temperature distribution,

$$T = T(x, y) = \underset{(1 \times 3)}{[N]} \underset{(3 \times 1)}{\{T\}}$$

$$\frac{\partial T}{\partial x} = \underset{(1 \times 3)}{\left[\frac{\partial N}{\partial x} \right]} \underset{(3 \times 1)}{\{T\}}$$

and

$$\frac{\partial T}{\partial y} = \underset{(1 \times 3)}{\left[\frac{\partial N}{\partial y} \right]} \underset{(3 \times 1)}{\{T\}}$$

Substitute to get the final form of element equations,

DERIVATION OF ELEMENT EQUATIONS

$$\int_{\Omega^{(e)}} \underbrace{\left(\begin{Bmatrix} \frac{\partial N}{\partial x} \\ \frac{\partial N}{\partial y} \end{Bmatrix}^k \begin{bmatrix} \frac{\partial N}{\partial x} \\ \frac{\partial N}{\partial y} \end{bmatrix} + \begin{Bmatrix} \frac{\partial N}{\partial y} \\ \frac{\partial N}{\partial x} \end{Bmatrix}^k \begin{bmatrix} \frac{\partial N}{\partial x} \\ \frac{\partial N}{\partial y} \end{bmatrix} \right)}_{\substack{[K_c] \\ (3 \times 3)}} d\Omega \quad \{T\}_{(3 \times 1)}$$

$$= \int_{\Gamma^{(e)}} \underbrace{\{N\}_{(3 \times 1)}}_{(3 \times 1)} \underbrace{\left(k \frac{\partial T}{\partial x} n_x + k \frac{\partial T}{\partial y} n_y \right)}_{\{Q_c\}_{(3 \times 1)}} d\Gamma + \int_{\Omega^{(e)}} \underbrace{\{N\}_{(3 \times 1)}}_{(3 \times 1)} Q d\Omega + \int_{s_2} \underbrace{\{N\}_{(3 \times 1)}}_{(3 \times 1)} q_s d\Gamma$$

Or, in short (details given in text):

$$[K_c] \{T\} = \{Q_c\} + \{Q_Q\} + \{Q_q\}$$

FINITE ELEMENT MATRICES

Element matrices for triangle can be derived in closed-form. Since,

$$N_i = \frac{1}{2A} (a_i + b_i x + c_i y) \quad i = 1, 2, 3$$

then $\frac{\partial N_i}{\partial x} = \frac{b_i}{2A}$ and $\frac{\partial N_i}{\partial y} = \frac{c_i}{2A}$

Conduction matrix Coefficients in the matrix are,

$$K_{ij} = \int_A k \left(\frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} + \frac{\partial N_i}{\partial y} \frac{\partial N_j}{\partial y} \right) t \, dx \, dy \quad i, j = 1, 2, 3$$

For constant thermal conductivity k and thickness t , then

$$K_{ij} = k t \int \left(\frac{b_i}{2A} \frac{b_j}{2A} + \frac{c_i}{2A} \frac{c_j}{2A} \right) dx \, dy$$

FINITE ELEMENT MATRICES

Or,
$$K_{ij} = \frac{k t}{4A} (b_i b_j + c_i c_j) \quad i, j = 1, 2, 3$$

$$\underset{(3 \times 3)}{[K_c]} = k A t \underset{(3 \times 2)}{[B]}^T \underset{(2 \times 3)}{[B]} \quad \text{where} \quad [B] = \frac{1}{2A} \begin{bmatrix} b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{bmatrix}$$

Load vector due to heat generation Coefficients are,

$$Q_i = \int_A N_i Q t \, dx \, dy \quad i = 1, 2, 3$$

For constant Q and t over element, integrate to get,

$$Q_i = \frac{Q A t}{3}$$

FINITE ELEMENT MATRICES

Or, in matrix form:
$$\{Q_Q\} = \frac{Q A t}{3} \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix}$$

Load vector due to edge heating

$$\{Q_q\} = \int_{s_2} \{N\} q_s \, d\Gamma \quad \text{E.g.,} \quad \{Q_q\} = \frac{q_s t \ell_{12}}{2} \begin{Bmatrix} 1 \\ 1 \\ 0 \end{Bmatrix}$$

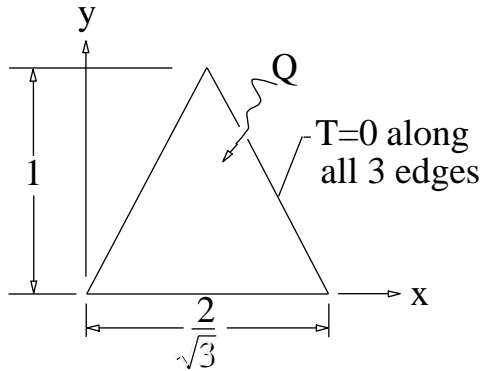
having nodes 1 and 2 on the edge.

Load vector for conduction

$$\{Q_c\} = \int_{\Gamma^{(e)}} \{N\} \left(k \frac{\partial \Gamma}{\partial x} n_x + k \frac{\partial \Gamma}{\partial y} n_y \right) d\Gamma$$

TWO-DIMENSIONAL HEAT TRANSFER

Example Determine temperature distribution in a triangular plate with internal heat generation using 3 finite elements.



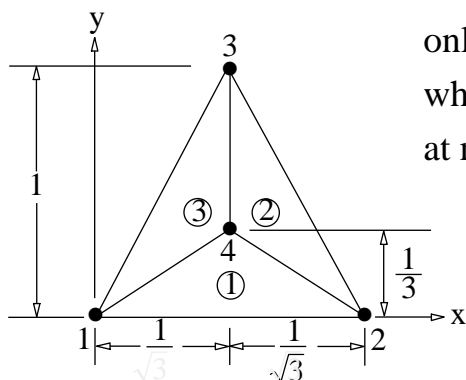
triangular plate with internal heat generation using 3 finite elements.

Note: This example has exact temperature solution as,

$$\bar{T}(x, y) = \frac{Q}{4k} (y - 2 + \sqrt{3}x)(y - \sqrt{3}x)y$$

TWO-DIMENSIONAL HEAT TRANSFER

we first discretize model into 3 elements. Here we have only one unknown at node 4 where as temperatures are zero at nodes 1, 2, and 3.



only one unknown at node 4 where as temperatures are zero at nodes 1, 2, and 3.

Typical element eqs. are:

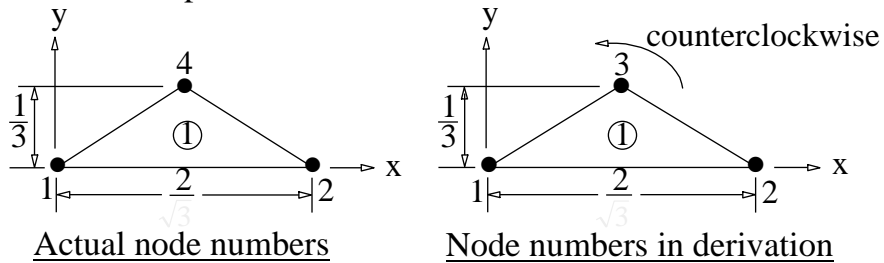
$$[K_c]\{T\} = \{Q_c\} + \{Q_Q\}$$

e.g., for element no. 1,

$$\begin{bmatrix} k_{11} & k_{12} & k_{13} \\ k_{21} & k_{22} & k_{23} \\ k_{31} & k_{32} & k_{33} \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \end{Bmatrix} = \begin{Bmatrix} Q_{c1} \\ Q_{c2} \\ Q_{c3} \end{Bmatrix} + \frac{QA_t}{3} \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix}$$

TWO-DIMENSIONAL HEAT TRANSFER

Computation of element matrices:



Here, $x_1 = 0$ $y_1 = 0$ $b_1 = -\frac{1}{3}$ $c_1 = -\frac{1}{\sqrt{3}}$
 $x_2 = \frac{2}{\sqrt{3}}$ $y_2 = 0$ $b_2 = \frac{1}{3}$ $c_2 = -\frac{1}{\sqrt{3}}$
 $x_3 = \frac{1}{\sqrt{3}}$ $y_3 = \frac{1}{3}$ $b_3 = 0$ $c_3 = \frac{2}{\sqrt{3}}$

with the element area $A = \frac{1}{3}\sqrt{3}$.

TWO-DIMENSIONAL HEAT TRANSFER

With these coefficients b_i and c_i , $i = 1, 2, 3$, element matrices can be determined. For example the coefficient K_{23} in the conduction matrix is,

$$\begin{aligned}
 K_{23} &= \frac{k t}{4A} (b_2 b_3 + c_2 c_3) \\
 &= \frac{k t}{4 \left(\frac{1}{3\sqrt{3}} \right)} \left(\left(\frac{1}{3} \right) (0) + \left(-\frac{1}{\sqrt{3}} \right) \left(\frac{2}{\sqrt{3}} \right) \right) = k t \left(-\frac{3}{2\sqrt{3}} \right)
 \end{aligned}$$

Thus, element equations for element no. ① are,

$$\frac{k t}{2\sqrt{3}} \begin{bmatrix} 2 & 1 & -3 \\ 1 & 2 & -3 \\ -3 & -3 & 6 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_4 \end{Bmatrix} = \begin{Bmatrix} (Q_{c_1})_1 \\ (Q_{c_2})_1 \\ (Q_{c_4})_1 \end{Bmatrix} + \frac{Q t}{9\sqrt{3}} \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix}$$

TWO-DIMENSIONAL HEAT TRANSFER

Similarly, element eqs. for element no. ② are

$$\frac{k t}{2\sqrt{3}} \begin{bmatrix} 2 & 1 & -3 \\ 1 & 2 & -3 \\ -3 & -3 & 6 \end{bmatrix} \begin{Bmatrix} T_2 \\ T_3 \\ T_4 \end{Bmatrix} = \begin{Bmatrix} (Q_{c_2})_2 \\ (Q_{c_3})_2 \\ (Q_{c_4})_2 \end{Bmatrix} + \frac{Q t}{9\sqrt{3}} \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix}$$

and for element no. ③ are,

$$\frac{k t}{2\sqrt{3}} \begin{bmatrix} 2 & 1 & -3 \\ 1 & 2 & -3 \\ -3 & -3 & 6 \end{bmatrix} \begin{Bmatrix} T_3 \\ T_1 \\ T_4 \end{Bmatrix} = \begin{Bmatrix} (Q_{c_3})_3 \\ (Q_{c_1})_3 \\ (Q_{c_4})_3 \end{Bmatrix} + \frac{Q t}{9\sqrt{3}} \begin{Bmatrix} 1 \\ 1 \\ 1 \end{Bmatrix}$$

Then assemble element equations. from these 3 elements to get,

TWO-DIMENSIONAL HEAT TRANSFER

$$\frac{k t}{2\sqrt{3}} \begin{bmatrix} 2+2 & 1 & 1 & -3-3 \\ 1 & 2+2 & 1 & -3-3 \\ 1 & 1 & 2+2 & -3-3 \\ -3-3 & -3-3 & -3-3 & 6+6+6 \end{bmatrix} \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \\ T_4 \end{Bmatrix} = \begin{Bmatrix} (Q_{c_1})_1 + (Q_{c_1})_3 \\ (Q_{c_2})_1 + (Q_{c_2})_2 \\ (Q_{c_3})_2 + (Q_{c_3})_3 \\ (Q_{c_4})_1 + (Q_{c_4})_2 + (Q_{c_4})_3 \end{Bmatrix} + \frac{Q t}{9\sqrt{3}} \begin{Bmatrix} 1+1 \\ 1+1 \\ 1+1 \\ 1+1+1 \end{Bmatrix}$$

TWO-DIMENSIONAL HEAT TRANSFER

Apply boundary conditions of $T_1 = T_2 = T_3 = 0$, the system eqs. become,

$$\frac{k t}{2\sqrt{3}} \begin{bmatrix} 4 & 1 & 1 & -6 \\ 1 & 4 & 1 & -6 \\ 1 & 1 & 4 & -6 \\ -6 & -6 & -6 & 13 \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ 0 \\ T_4 \end{Bmatrix} = \begin{Bmatrix} Q_1 \\ Q_2 \\ Q_3 \\ 0 \end{Bmatrix} + \frac{Q t}{9\sqrt{3}} \begin{Bmatrix} 2 \\ 2 \\ 2 \\ 3 \end{Bmatrix}$$

Then use the last equation to solve for temperature at node 4,

$$\frac{k t}{2\sqrt{3}} (0 + 0 + 0 + 13 T_4) = 0 + \frac{Q t}{9\sqrt{3}} (3)$$

$$T_4 = \frac{1}{27} \frac{Q}{k}$$

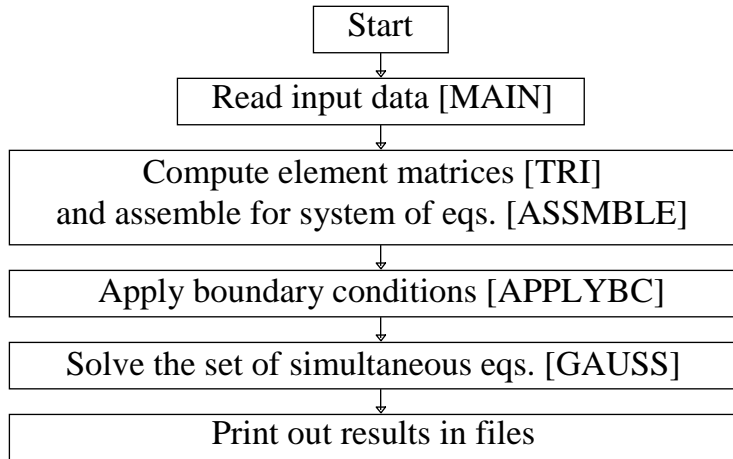
and then can use the first 3 equations to solve for heat fluxes at nodes,

$$Q_1 = Q_2 = Q_3 = -\frac{6 k t}{2\sqrt{3}} T_4 - \frac{2 Q t}{9\sqrt{3}} = -\frac{Q t}{3\sqrt{3}}$$

COMPUTER PROGRAMMING

- Same procedure described in example can be used in the development of finite element computer program directly.
- Listing of sample finite element program provided in following pages (as short as only 6 A-4 pages).
- Program is applicable to arbitrary geometry.
- Program can be modified to solve other problems.
- By understanding this short program -- won't be afraid of "black box" anymore.

COMPUTER PROGRAMMING FLOW-CHART



Note: Names in the square brackets [] represent the subroutine names in the finite element program illustrated.

FINITE ELEMENT PROGRAM LISTING ²³⁰

Example in FORTRAN

```

C      PROGRAM FINITE
C
C      A FINITE ELEMENT COMPUTER PROGRAM FOR SOLVING PARTIAL
C      DIFFERENTIAL EQUATION IN THE FORM OF POISSON'S EQUATION
C      FOR TWO-DIMENSIONAL STEADY-STATE HEAT CONDUCTION WITH
C      INTERNAL HEAT GENERATION.
C
C      DR. PRAMOTE DECHADUMPHAI
C      FACULTY OF ENGINEERING
C      CHEULALONGKORN UNIVERSITY
C
C      THE VALUES DECLARED IN THE PARAMETER STATEMENT BELOW SHOULD
C      BE ADJUSTED ACCORDING TO THE SIZE OF THE PROBLEMS AND TYPES
C      OF COMPUTERS.
C
C      MKPOI = MAXIMUM NUMBER OF NODES IN THE MODEL
C      MKELE = MAXIMUM NUMBER OF ELEMENTS IN THE MODEL
C
C      PARAMETER (MKPOI=150, MKELE=300)
C
C      IMPLICIT REAL*8 (A-H,O-Z)
C      DIMENSION COORD(MKPOI,2), TEMP(MKPOI), TEXT(20)
C      DIMENSION SYEQ(MKPOI,MKPOI), SYEQ(MKPOI), QELB(MKLE)
C      CHARACTER*20 NAME1, NAME2
C
C      INTEGER INTMAT(MKLE,3), INC(MKPOI)
C
C      10 WRITE(6,20)
C      20 FORMAT(' ', ' PLEASE ENTER THE INPUT FILE NAME:')
C      READ(5, '(A)', ERR=10) NAME1
C      OPEN(UNIT=7, FILE=NAME1, STATUS='OLD', ERR=10)
C
C      C      READ TITLE OF COMPUTATION:
C      READ(7,*) NLINE
C      DO 100 ILINE=1,NLINE
C      READ(7,1) TEXT
C      1 FORMAT(20A4)
C      100 CONTINUE
C
C      C      READ INPUT DATA:
C      READ(7,1) TEXT
C      READ(7,*) NPQIN, NLELM
C      IF(FUNCTION.GT.MKPOI) WRITE(6,110) NPQIN
C      110 FORMAT(' ', ' PLEASE INCREASE THE PARAMETER MKPOI TO ', 15)
C      IF(FUNCTION.GT.MKLE) WRITE(6,120) NLELM
C      120 FORMAT(' ', ' PLEASE INCREASE THE PARAMETER MKLE TO ', 15)
C      IF(FUNCTION.GT.MKLE) STOP
C      READ(7,1) TEXT
C      READ(7,*) TX, THICK
C      READ(7,1) TEXT
C      DO 130 I=1,NPQIN
C      READ(7,*) I, INC(I), (COORD(I,K), K=1,2), TEMP(I)
C      IF(I.NE.IP) WRITE(6,135) IP
C      135 FORMAT(' ', ' NODE NO. ', 15, ' IN DATA FILE IS MISSING')
C      IF(I.NE.IP) STOP
C      130 CONTINUE
C      140 = 0
C      READ(7,1) TEXT
C      DO 140 I=1,NLELM
C
C      READ(7,*) I, (INTMAT(I,J), J=1,3), QELB(I)
C      IF(I.NE.IE) WRITE(6,150) IE
C      150 FORMAT(' ', ' ELEMENT NO. ', 15, ' IN DATA FILE IS MISSING')
C      IF(I.NE.IE) STOP
C      IF(QELB(I).NE.0.) IQ = 1
C      140 CONTINUE
C      WRITE(6,160)
C      160 FORMAT(' ', ' THE F.E. MODEL INCLUDES THE FOLLOWING',
C      *      ' HEAT TRANSFER MODE(S):',
C      *      ' ', ' -- HEAT CONDUCTION ',
C      *      ' IF(IQ.EQ.1) WRITE(6,170) ',
C      *      ' 170 FORMAT(' ', ' -- INTERNAL HEAT GENERATION ')
C
C      NRQ = NPQIN
C      DO 180 I=1,NRQ
C      SYEQ(I) = 0.
C      180 CONTINUE
C      DO 190 J=1,NRQ
C      DO 190 I=1,NRQ
C      SYEQ(I,J) = 0.
C      190 CONTINUE
C
C      WRITE(6,200) NPQIN, NLELM
C      200 FORMAT(' ', ' *** THE FINITE ELEMENT MODEL CONSISTS OF ', 15,
C      *      ' NODES AND ', 15, ' ELEMENTS ***')
C
C      C      ESTABLISH ALL ELEMENT MATRICES AND ASSEMBLE THEM TO FORM
C      C      FORM UP SYSTEM EQUATIONS
C      WRITE(6,210)
C      210 FORMAT(' ', ' *** ESTABLISHING ELEMENT MATRICES AND',
C      *      ' ASSEMBLING ELEMENT EQUATIONS ***')
C      CALL TRN(NLELM, INTMAT, COORD, TX, QELB, THICK,
C      *      SYEQ, SYEQ, MKPOI, MKLE)
C
C      WRITE(6,220)
C      220 FORMAT(' ', ' *** APPLYING BOUNDARY CONDITIONS OF NODAL',
C      *      ' TEMPERATURES ***')
C      CALL APPLYBC(NPQIN, INC, TEMP, SYEQ, SYEQ, MKPOI)
C
C      WRITE(6,230)
C      230 FORMAT(' ', ' *** SOLVING A SET OF SIMULTANEOUS EQUATIONS',
C      *      ' FOR TEMPERATURE SOLUTIONS ***')
C      WRITE(6,240) NRQ
C      240 FORMAT(15, '( SOLVE OF ', 15, ' EQUATIONS TO BE SOLVED )')
C      CALL GAUSS(NRQ, SYEQ, SYEQ, TEMP, MKPOI)
C
C      C      PRINT OUT NODAL TEMPERATURE SOLUTIONS:
C      C
C      250 WRITE(6,260)
C      260 FORMAT(' ', ' PLEASE ENTER FILE NAME FOR TEMPERATURE',
C      *      ' SOLUTIONS:')
C      READ(5, '(A)', ERR=250) NAME2
C      OPEN(UNIT=8, FILE=NAME2, STATUS='NEW', ERR=250)
C      WRITE(8,270) NPQIN
C      270 FORMAT(' ', ' NODAL TEMPERATURE SOLUTIONS [ ', 15, ']:',
C      *      ' //, 2X, ' NODE', 3X, ' TEMPERATURE', /
C      *      ' DO 280 I=1,NPQIN
C      *      WRITE(8,290) IP, TEMP(IP)
C      *      280 FORMAT(15, ' E14.6)
C      280 CONTINUE
C
C      STOP
C      END
  
```

FINITE ELEMENT PROGRAM LISTING

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C
C SUBROUTINE APPLVBC(NP0IN, IBC, TEMP, SYEQ, SYSQ, MKPOI)
C
C APPLY TEMPERATURE BOUNDARY CONDITIONS WITH CONDITION CODES OF:
C 0 = FREE TO CHANGE (TO BE COMPUTED)
C 1 = FIXED AS SPECIFIED
C
C
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION SYEQ(MKPOI,MKPOI), SYSQ(MKPOI), TEMP(MKPOI)
C
C INTEGER IBC(MKPOI)
C
C DO 100 IBC=1,NP0IN
C IF(IBC(IBC),EQ.0) GO TO 100
C
C DO 200 IBC=1,NP0IN
C IF(IBC(IBC),EQ.0) GO TO 200
C SYEQ(IBC) = SYSQ(IBC) - SYEQ(IBC)*TEMP(IBC)
C SYEQ(IBC) = 0.
C 200 CONTINUE
C
C DO 300 IBC=1,NP0IN
C SYEQ(IBC,IBC) = 0.
C 300 CONTINUE
C SYEQ(IBC,IBC) = 1.
C SYEQ(IBC) = TEMP(IBC)
C
C 100 CONTINUE
C RETURN
C END
C
C-----
C SUBROUTINE ASSEMBLE( IE, INTMAT, AKC, QQ, SYEQ, SYSQ,
C MKPOI, MKELE )
C
C ASSEMBLE ELEMENT EQUATIONS INTO SYSTEM EQUATIONS
C
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION AKC(3,3), QQ(3)
C DIMENSION SYEQ(MKPOI,MKPOI), SYSQ(MKPOI)
C
C INTEGER INTMAT(MKLE,3)
C
C MKLE = 3
C
C DO 100 IBC=1,NKODE
C DO 200 IBC=1,NKODE
C IBCW = INTMAT(IBC,1)
C ICOL = INTMAT(IBC,2)
C SYEQ(IBCW,ICOL) = SYEQ(IBCW,ICOL) + AKC(IBC,IBC)
C 200 CONTINUE
C SYEQ(IBCW) = SYEQ(IBCW) + QQ(IBC)
C 100 CONTINUE
C
C RETURN
C END
C
C-----
C SUBROUTINE GAUSS(N, A, B, X, MKPOI)
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION A(MKPOI,MKPOI), B(MKPOI), X(MKPOI)
C
C PERFORM SCALING:
C CALL SCALE(N, A, B, MKPOI)

```

```

C
C FORWARD ELIMINATION:
C
C PERFORM ACCORDING TO ORDER OF 'PRIME' FROM 1 TO N-1:
C
C DO 100 I=1,N-1
C
C PERFORM PARTIAL PIVOTING:
C CALL PIVOT(N, A, B, MKPOI, IP)
C
C LOOP OVER EACH EQUATION STARTING FROM THE ONE THAT CORRESPONDS
C WITH THE ORDER OF 'PRIME' PLUS ONE:
C
C DO 200 IE=IP+1,N
C RATIO = A(IE,IP)/A(IP,IP)
C
C COMPUTE NEW COEFFICIENTS OF THE EQUATION CONSIDERED:
C
C DO 300 IC=IP+1,N
C A(IE,IC) = A(IE,IC) - RATIO*A(IP,IC)
C 300 CONTINUE
C B(IE) = B(IE) - RATIO*B(IP)
C 200 CONTINUE
C
C SET COEFFICIENTS ON LOWER LEFT PORTION TO ZERO:
C
C DO 400 IE=IP+1,N
C A(IE,IP) = 0.
C 400 CONTINUE
C 100 CONTINUE
C
C BACK SUBSTITUTION:
C
C COMPUTE SOLUTION OF THE LAST EQUATION:
C X(N) = B(N)/A(N,N)
C
C THEN COMPUTE SOLUTIONS FROM EQUATION N-1 TO 1:
C
C DO 500 IE=N-1,1,-1
C SUM = 0.
C DO 600 IC=IE+1,N
C SUM = SUM + A(IE,IC)*X(IC)
C 600 CONTINUE
C X(IE) = (B(IE) - SUM)/A(IE,IE)
C 500 CONTINUE
C RETURN
C END
C
C-----
C SUBROUTINE PIVOT(N, A, B, MKPOI, IP)
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION A(MKPOI,MKPOI), B(MKPOI)
C
C PERFORM PARTIAL PIVOTING:
C
C JP = IP
C BIG = ABS(A(IP,IP))
C DO 10 I=IP+1,N
C AMAX = ABS(A(I,IP))
C IF(AMAX.GT.BIG) THEN
C BIG = AMAX
C JP = I
C ENDIF
C 10 CONTINUE

```

FINITE ELEMENT PROGRAM LISTING

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C
C IF(AMAX.GT.BIG) THEN
C BIG = AMAX
C JP = I
C ENDDIF
C 10 CONTINUE
C IF(JP.NE.IP) THEN
C DINT = A(IP,JP)
C A(IP,JP) = A(JP,JP)
C A(JP,JP) = DINT
C A(IP,JP) = DINT
C 20 CONTINUE
C DINT = B(JP)
C B(IP) = B(IP)
C B(JP) = DINT
C ENDDIF
C RETURN
C END
C
C-----
C SUBROUTINE SCALE(N, A, B, MKPOI)
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION A(MKPOI,MKPOI), B(MKPOI)
C
C PERFORM SCALING:
C
C DO 10 I=1,N
C BIG = ABS(A(IE,1))
C DO 20 IC=2,N
C AMAX = ABS(A(IE,IC))
C IF(AMAX.GT.BIG) BIG = AMAX
C 20 CONTINUE
C DO 30 IC=1,N
C A(IE,IC) = A(IE,IC)/BIG
C 30 CONTINUE
C B(IE) = B(IE)/BIG
C 10 CONTINUE
C RETURN
C END
C
C-----
C SUBROUTINE TRI(INTELM, INTMAT, COORD, TK, QELE, THICK,
C SYEQ, SYSQ, MKPOI, MKELE )
C
C ESTABLISH ALL ELEMENT MATRICES AND ASSEMBLE THEM TO FORM
C UP SYSTEM EQUATIONS
C
C IMPLICIT REAL*8 (A-H,O-Z)
C DIMENSION COORD(MKPOI,2), SYEQ(MKPOI,MKPOI), SYSQ(MKPOI)
C DIMENSION QELE(MKLE)
C DIMENSION AKC(3,3), QQ(3), B(2,3), BT(3,2)
C
C INTEGER INTMAT(MKLE,3)
C
C LOOP OVER THE NUMBER OF ELEMENTS:
C
C DO 500 I=1,NELEM
C
C FIND ELEMENT LOCAL COORDINATES:
C
C II = INTMAT(IE,1)
C JJ = INTMAT(IE,2)
C KK = INTMAT(IE,3)
C
C X01 = COORD(II,1)
C X02 = COORD(JJ,1)
C X03 = COORD(KK,1)

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```

C
C Y01 = COORD(II,2)
C Y02 = COORD(JJ,2)
C Y03 = COORD(KK,2)
C
C AREA = 0.5*(X02*(Y03-Y01) + X01*(Y02-Y03) + X03*(Y01-Y02))
C IF(AREA.LE.0.) WRITE(6,5) IE
C 5 FORMAT(' ', I11, ' ERROR !!! ELEMENT NO. ', I5,
C ' HAS NEGATIVE OR ZERO AREA ', /,
C ' --- CHECK F.E. MODEL FOR NODAL COORDINATES',
C ' AND ELEMENT NODAL CONNECTIONS ---')
C IF(AREA.LE.0.) STOP
C
C B1 = Y02 - Y03
C B2 = Y03 - Y01
C B3 = Y01 - Y02
C C1 = X02 - X03
C C2 = X03 - X01
C C3 = X01 - X02
C
C DO 10 I=1,2
C DO 10 J=1,2
C B(I,J) = B(I,J)
C 10 CONTINUE
C
C B(1,1) = B1
C B(1,2) = B2
C B(1,3) = B3
C B(2,1) = C1
C B(2,2) = C2
C B(2,3) = C3
C
C DO 10 I=1,2
C DO 10 J=1,3
C B(I,J) = B(I,J)/(2.*AREA)
C BT(I,J) = B(I,J)
C 10 CONTINUE
C
C ELEMENT CONDUCTION MATRIX:
C
C DO 100 I=1,3
C DO 100 J=1,3
C AKC(I,J) = 0.
C DO 110 IK=1,2
C AKC(I,J) = AKC(I,J) + BT(I,K)*B(K,J)
C 110 CONTINUE
C AKC(I,J) = TK*AREA*THICK*AKC(I,J)
C 100 CONTINUE
C
C ELEMENT LOAD VECTOR DUE TO INTERNAL HEAT GENERATION:
C
C FAC = QELE(IE)*AREA*THICK/3.
C DO 200 I=1,3
C QQ(I) = FAC
C 200 CONTINUE
C
C ASSEMBLE THESE ELEMENT MATRICES TO FORM SYSTEM EQUATIONS:
C CALL ASSEMBLE( IE, INTMAT, AKC, QQ, SYEQ, SYSQ,
C MKPOI, MKELE )
C
C 500 CONTINUE
C RETURN
C END

```

TYPICAL INPUT & OUTPUT DATA

Input data

```

2
TRIANGULAR PLATE WITH INTERNAL HEAT GENERATION.
CRUDE MESH WITH 4 NODES AND 3 ELEMENTS.
NPOIN  NELEM
 4      3
TK      THICK
1.      .1
NODAL BOUNDARY CONDITIONS AND COORDINATES [4]:
 1      1      0.00000  0.00000  0.
 2      1      1.15470  0.00000  0.
 3      1      0.57735  1.00000  0.
 4      0      0.57735  0.33333  0.
ELEMENT NODAL CONNECTIONS AND HEAT GEN. [3]:
 1      1      2      4      1.
 2      2      3      4      1.
 3      3      1      4      1.

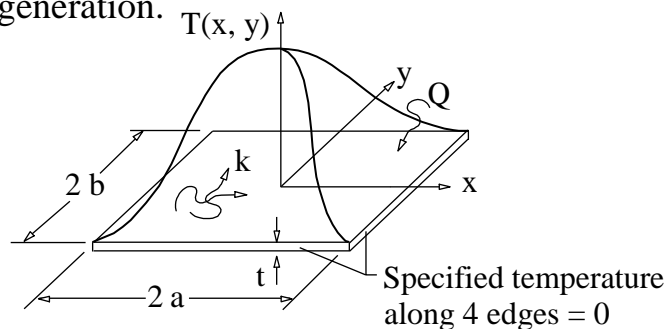
```

Output result

NODE	TEMPERATURE
1	.000000E+00
2	.000000E+00
3	.000000E+00
4	.370370E-01

RECTANGULAR ELEMENT

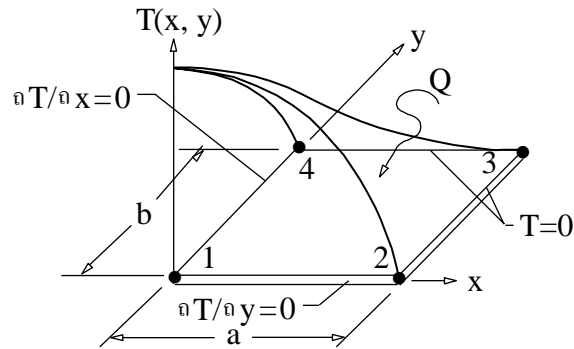
Example Two-dimensional steady-state heat conduction in a rectangular plate with internal heat generation. $T(x, y)$



Due to symmetry, only a quarter of the plate can be modeled. Use one 4-node rectangular element to compute temperature at the plate center ($x=y=0$).

RECTANGULAR ELEMENT

By using only one rectangular element for quarter of the plate, some edge conditions must be provided,

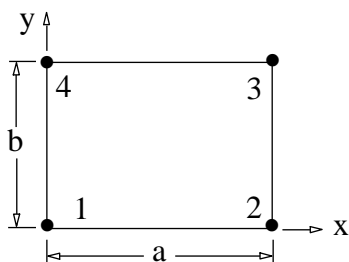


Here
Want

$$\begin{aligned} T_2 &= T_3 = T_4 = 0 \\ T_1 &= ? \end{aligned}$$

RECTANGULAR ELEMENT

Element interpolation function:



Assume element temperature,
 $T(x, y) = \alpha_1 + \alpha_2 x + \alpha_3 y + \alpha_4 xy$
 which is in form of bilinear distribution where α_i , $i = 1, 4$ are constants that can be determined from the conditions,

$$\begin{aligned} T(0,0) &= T_1 = \alpha_1 + 0 + 0 + 0 \\ T(a,0) &= T_2 = \alpha_1 + \alpha_2 a + 0 + 0 \\ T(a,b) &= T_3 = \alpha_1 + \alpha_2 a + \alpha_3 b + \alpha_4 ab \\ T(0,b) &= T_4 = \alpha_1 + 0 + \alpha_3 b + 0 \end{aligned}$$

RECTANGULAR ELEMENT

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Solve for $\alpha_i, i = 1, 4$ to get,

$$\begin{aligned} \alpha_1 &= T_1 & \alpha_3 &= \frac{T_4 - T_1}{b} \\ \alpha_2 &= \frac{T_2 - T_1}{a} & \alpha_4 &= \frac{T_1 - T_2 + T_3 - T_4}{ab} \end{aligned}$$

Then, substitute back and rearrange terms to get element temperature distribution in the form,

$$\begin{aligned} T &= \left[\underbrace{\left(1 - \frac{x}{a}\right) \left(1 - \frac{y}{b}\right)}_{N_1} \quad \underbrace{\frac{x}{a} \left(1 - \frac{y}{b}\right)}_{N_2} \quad \underbrace{\frac{x}{a} \frac{y}{b}}_{N_3} \quad \underbrace{\frac{y}{b} \left(1 - \frac{x}{a}\right)}_{N_4} \right] \begin{Bmatrix} T_1 \\ T_2 \\ T_3 \\ T_4 \end{Bmatrix} \\ &= \underset{(1 \times 4)}{[N(x, y)]} \underset{(4 \times 1)}{\{T\}} \end{aligned}$$

where $[N(x, y)]$ is the element interpolation function matrix.

PLATE WITH INTERNAL HEAT GEN.

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Since the corresponding FE eqs. for this problem are,

$$[K_c] \{T\} = \{Q_c\} + \{Q_Q\}$$

where $[K_c]$ is the conduction matrix,

$$[K_c] = \int_{\Omega^{(e)}} k \left(\left\{ \frac{\partial N}{\partial x} \right\} \left[\frac{\partial N}{\partial x} \right] + \left\{ \frac{\partial N}{\partial y} \right\} \left[\frac{\partial N}{\partial y} \right] \right) d\Omega$$

Here,

$$\begin{aligned} \left[\frac{\partial N}{\partial x} \right] &= \left[-\frac{1}{a} \left(1 - \frac{y}{b}\right) \quad \frac{1}{a} \left(1 - \frac{y}{b}\right) \quad \frac{1}{a} \frac{y}{b} \quad -\frac{1}{a} \frac{y}{b} \right] \\ \left[\frac{\partial N}{\partial y} \right] &= \left[-\frac{1}{b} \left(1 - \frac{x}{a}\right) \quad -\frac{1}{b} \frac{x}{a} \quad \frac{1}{b} \frac{x}{a} \quad \frac{1}{b} \left(1 - \frac{x}{a}\right) \right] \end{aligned}$$

and $d\Omega = t \, dx \, dy$. Perform integration with the limits $x = 0 \rightarrow a$ and $y = 0 \rightarrow b$ to obtain,

FINITE ELEMENT EQUATIONS

Conduction matrix,

$$[K_c] = k t \begin{bmatrix} \frac{1}{3} \left(\frac{a}{b} + \frac{b}{a} \right) & \frac{1}{6} \left(\frac{a}{b} - \frac{2b}{a} \right) & \frac{1}{6} \left(-\frac{a}{b} - \frac{b}{a} \right) & \frac{1}{6} \left(-\frac{2a}{b} + \frac{b}{a} \right) \\ & \frac{1}{3} \left(\frac{a}{b} + \frac{b}{a} \right) & \frac{1}{6} \left(-\frac{2a}{b} + \frac{b}{a} \right) & \frac{1}{6} \left(-\frac{a}{b} - \frac{b}{a} \right) \\ & & \frac{1}{3} \left(\frac{a}{b} + \frac{b}{a} \right) & \frac{1}{6} \left(\frac{a}{b} - \frac{2b}{a} \right) \\ \text{sym} & & & \frac{1}{3} \left(\frac{a}{b} + \frac{b}{a} \right) \end{bmatrix}$$

Also, the load vector from internal heat generation,

$$\{Q_Q\} = \int_{\Omega^{(e)}} Q \{N\} d\Omega = Q a b t \begin{Bmatrix} \frac{1}{4} \\ \frac{1}{4} \\ \frac{1}{4} \\ \frac{1}{4} \end{Bmatrix}$$

FINITE ELEMENT EQUATIONS

and the conduction vector,

$$\{Q_c\} = \int_{\Gamma^{(e)}} \{N\} k \left(\frac{\partial \Gamma}{\partial x} n_x + \frac{\partial \Gamma}{\partial y} n_y \right) d\Gamma = \begin{Bmatrix} Q_{c_1} \\ Q_{c_2} \\ Q_{c_3} \\ Q_{c_4} \end{Bmatrix}$$

Therefore, the final element eqs. can be written as,

$$\begin{bmatrix} K_{11} & K_{12} & K_{13} & K_{14} \\ & K_{22} & K_{23} & K_{24} \\ & & K_{33} & K_{34} \\ \text{Sym} & & & K_{44} \end{bmatrix} \begin{Bmatrix} T_1 = ? \\ T_2 = 0 \\ T_3 = 0 \\ T_4 = 0 \end{Bmatrix} = \begin{Bmatrix} Q_{c_1} \\ Q_{c_2} \\ Q_{c_3} \\ Q_{c_4} \end{Bmatrix} + Q a b t \begin{Bmatrix} \frac{1}{4} \\ \frac{1}{4} \\ \frac{1}{4} \\ \frac{1}{4} \end{Bmatrix}$$

Then use the first eq. to solve for T_1 ,

PLATE WITH INTERNAL HEAT GEN.

$$K_{11}T_1 + 0 + 0 + 0 = Q_{c_1} + \frac{Q a b t}{4}$$

$$\frac{k t}{3} \left(\frac{a}{b} + \frac{b}{a} \right) T_1 = Q_{c_1} + \frac{Q a b t}{4}$$

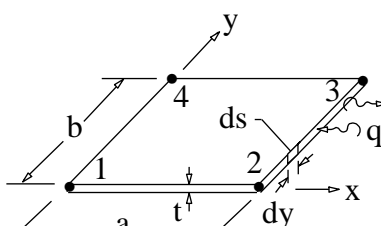
It can be shown that $Q_{c_1} = 0$ (no conduction across center of the plate), thus the temperature at the plate center is,

$$T_1 = \frac{3}{4} \frac{Q a b}{k} \frac{1}{\left(\frac{a}{b} + \frac{b}{a} \right)}$$

Note that if $Q = k = a = b = 1$, the exact solution is 0.295 where T_1 is 0.375. But if use 4 elements, T_1 becomes 0.315, i.e., solution is improved with more elements.

PLATE WITH EDGE HEAT TRANSFER

If the plate has specified heating and convection along edge 2-3. Since the load vector due to edge heating is,



$$\{Q_q\} = \int_{s_2}^{s_3} q_s \{N\} ds$$

along edge $x = a$, $ds = t dy$ and

$$[N(x = a, y)] = \begin{bmatrix} 0 & 1 - \frac{y}{b} & \frac{y}{b} & 0 \end{bmatrix}$$

Then,

$$\{Q_q\} = \int_0^b q_s \begin{Bmatrix} 0 \\ 1 - \frac{y}{b} \\ \frac{y}{b} \\ 0 \end{Bmatrix} t dy = q_s b t \begin{Bmatrix} 0 \\ \frac{1}{2} \\ \frac{1}{2} \\ 0 \end{Bmatrix}$$

PLATE WITH EDGE HEAT TRANSFER

Similarly, the convection matrix corresponding to convection heat transfer along edge 2-3 can be derived,

$$\begin{aligned}
 [K_h] &= \int_{s_2^{(e)}} h \{N\} [N] ds \\
 &= \int_0^b h \begin{Bmatrix} 0 \\ 1 - \frac{y}{b} \\ \frac{y}{b} \\ 0 \end{Bmatrix} \begin{bmatrix} 0 & 1 - \frac{y}{b} & \frac{y}{b} & 0 \end{bmatrix} t dy = h b t \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & \frac{1}{3} & \frac{1}{6} & 0 \\ 0 & \frac{1}{6} & \frac{1}{3} & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}
 \end{aligned}$$

PLATE WITH EDGE HEAT TRANSFER

with the convection load vector,

$$\begin{aligned}
 \{Q_h\} &= \int_{s_2^{(e)}} h T_\infty \{N\} ds = \int_{s_2^{(e)}} h T_\infty \begin{Bmatrix} 0 \\ 1 - \frac{y}{b} \\ \frac{y}{b} \\ 0 \end{Bmatrix} t dy \\
 &= h b t T_\infty \begin{Bmatrix} 0 \\ \frac{1}{2} \\ \frac{1}{2} \\ 0 \end{Bmatrix}
 \end{aligned}$$